

QUESTION 2: [40 points]

a) Consider the operator

$$T(a) = \exp(-iaP/\hbar) \quad (0.9)$$

Where  $a$  is a constant and  $P$  is the momentum operator Show that

$$T^\dagger(a) X T(a) = x + a \quad (0.10)$$

$$[X, P] = i\hbar \quad \text{use} \quad e^{+A} B e^{-A} = \exp \text{Ad}_A B$$

$$\begin{aligned} e^{+iaP/\hbar} X e^{-iaP/\hbar} &= (1 + \text{ad}_{iaP/\hbar}) X \\ &= X + i\frac{a}{\hbar} [P, X] \\ &= X + a \end{aligned}$$

b) Show that  $T(a)$  is unitary and show that  $T(a)$  has eigenvalues of the form  $e^{i\phi}$  where  $\phi$  is real (You can assume that  $P$  is hermitian).

$$T(a)^\dagger T(a) = e^{iaP/\hbar} e^{-iaP/\hbar} = 1 \quad \text{since } P \text{ commutes with itself}$$

$$\text{same for } T(a) T(a)^\dagger \Rightarrow T \text{ unitary}$$

$$\langle \psi | T(a) | \psi \rangle = \langle \psi | \lambda | \psi \rangle = \lambda$$

$$\langle \psi | T(a) | \psi \rangle^* = \langle \psi | T(a)^\dagger | \psi \rangle = \lambda^*$$

on the other hand

$$\begin{aligned} | \psi \rangle &= \lambda T(a)^\dagger | \psi \rangle \Rightarrow \lambda^* = \frac{1}{\lambda} \quad \lambda \lambda^* = 1 \\ &\Rightarrow \lambda = e^{i\phi} \end{aligned}$$

c) Consider the Hamiltonian which is periodic under shifts by  $a$

post quanta

$$H = \frac{P^2}{2m} + \sum_{n=-\infty}^{\infty} V(X - na) \quad (0.11)$$

Here you can assume that  $V(x)$  goes exponentially fast to zero as  $|x| \rightarrow \infty$  (This assumption makes the sum over  $n$  convergent). You can also assume that  $V(x)$  can be expanded in a power series.

Prove that  $T(a)$  commutes with  $H$ .

$$V(x) = \sum_n c_n x^n$$

$$T^+ X T = X + a$$

$$\begin{aligned} T^+ X^n T &= T^+ X T T^+ \dots T T^+ X T \\ &= (X + a)^n \end{aligned}$$

$$\begin{aligned} \Rightarrow T^+ V(x) T &= \sum_n c_n (X + a)^n \\ &= V(x + a) \end{aligned}$$

$$\begin{aligned} T^+ \sum_{n=-\infty}^{\infty} V(X - na) T &= \sum_{n=-\infty}^{\infty} V(X + a - na) \\ &= \sum_n V(X - (n-1)a) \end{aligned}$$

$$= \sum_{n'} V(X - n'a) = \sum V(X - na)$$

$$T^+ P^2 T = P^2$$

we

$$T^+ H T = H$$

$$HT = TH \Leftrightarrow [T, H] = 0$$

d) It follows from the results in part c) that the Hamiltonian  $H$  and  $T(a)$  can be diagonalized simultaneously. You can assume that there are eigenstates  $|E, k\rangle$  which satisfy

$$H |E, k\rangle = E |E, k\rangle \quad (0.12)$$

$$T(a) |E, k\rangle = e^{-ika} |E, k\rangle \quad (0.13)$$

For the wave functions in position space define the following combination

$$u_k(x) = \langle x | E, k \rangle e^{-ikx} \quad (0.14)$$

Show that  $u_k(x)$  is a periodic function with period  $a$ , i.e.

$$u_k(x+a) = u_k(x) \quad (0.15)$$

This is the Bloch's theorem for periodic potentials (i.e. an energy eigenstate can be written as a Bloch wave times a periodic function).

$$\begin{aligned} u_k(x+a) &= \langle x+a | E, k \rangle e^{-ik(x+a)} \\ &= (\langle E, k | x+a \rangle)^* e^{-ik(x+a)} \\ &= \langle E, k | T | x \rangle^* e^{-ik(x+a)} \\ &= \langle x | T^\dagger | E, k \rangle e^{-ikx} e^{-ika} \\ &= \langle x | E, k \rangle e^{ika} e^{-ikx} e^{-ika} \\ &= u_k(x) \end{aligned}$$

Use

$$T_{(a)}^\dagger T_{(a)} |E, k\rangle = e^{-ika} T^\dagger |E, k\rangle$$

$$\Rightarrow T^\dagger |E, k\rangle = e^{ika} |E, k\rangle$$